



Fig. 3 Comparison of  $\beta^2$  of the first three modes in a shield microstrip line between accurate and approximate data. Parameters:  $h_1 = 3$  mm,  $h_2 = h_3 = 0.3175$  mm,  $w = 0.56$  mm,  $A = 5$  mm,  $\epsilon_{r3} = 10$ ,  $\sigma_2 = \sigma_3 = 0$ . Basic functions: ( $\beta_1 = 0.28709E-01$ , 0.0,  $\text{Freq}_1 = 1.0$  GHz). Near cutoff frequency: ( $\beta_2 = 1.0644E-001$ , 0.0,  $\text{Freq}_2 = 16.04$  GHz), ( $\beta_3 = 1.9356E-001$ , 0.0,  $\text{Freq}_3 = 29.20$  GHz).

#### REFERENCES

- [1] P. Przybyszewski, J. Mielewski, and Mrozowski, "Fast technique for analysis of waveguides," *IEEE Microwave Guided Wave Lett.*, vol. 8, pp. 109–111, Mar. 1998.
- [2] P. Przybyszewski, J. Mielewski, and M. Mrozowski, "A new class of eigenfunction expansion methods for fast frequency-domain analysis of waveguides," *IEEE Trans. Microwave Theory Tech.*, vol. 51, pp. 558–563, Feb. 2002.

### Comments on "Rigorous Modeling of Packaged Schottky Diodes by the Nonlinear Lumped Network (NL<sup>2</sup>N)–FDTD Approach"

Otman El Mrabet and Mohamed Essaaidi

In the above paper,<sup>1</sup> we have found several errors in (1), (3), (4), (6), (9), and (14), and those of the parameters presented in Tables I and II that should be corrected to allow an accurate implementation

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The authors are with the Electronics and Microwaves Group, Abdelmalek Essaadi University, Tetuan 93000, Morocco.

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<sup>1</sup>G. Emili, F. Alimenti, P. Mezzanote, L. Roselli, and R. Sorrentino, *IEEE Trans. Microwave Theory Tech.*, vol. 48, no. 12, pp. 2277–2282, Dec. 2000.

of the model proposed in the above paper for the modeling of packaged Schottky diodes by the nonlinear lumped-network finite-difference time-domain (FDTD) approach.

Considering the Schottky diode model depicted in Fig. 1(b) in the above paper, we can easily prove that the inductance  $L_1$  should be replaced by  $L_2$  in the expression of the admittance matrix of this circuit given by (1). Consequently, this admittance matrix should be written as follows:

$$[Y] = \frac{1}{D(s)} \begin{bmatrix} L_1 C s^2 + R C' s + 1 & -1 \\ -1 & L_2 C s^2 + 1 \end{bmatrix}$$

$$D(s) = L_1 L_2 C s^3 + R L_2 C s^2 + (L_1 + L_2) s + R. \quad (1)$$

Equation (3) should be written without the term "1" appearing in the denominator as follows:

$$\hat{Y}_{pq}(z) = \frac{\sum_{r=0}^M C_r^{pq} Z^{-r}}{\sum_{r=0}^M d_r Z^{-r}}. \quad (2)$$

The left-hand side of all the expressions in (4) and (6) of the above paper should be multiplied by  $d_0$ . Therefore, the correct form of these equations should be written, respectively, as follows:

$$d_0 I_1^{n+1} = \sum_{r=0}^3 C_r^{11} V_1^{n-r+1} + \sum_{r=0}^3 C_r^{12} V_2^{n-r+1} - \sum_{r=1}^3 d_r I_1^{n-r+1} \quad (3)$$

$$d_0 I_2^{n+1} = \sum_{r=0}^3 C_r^{21} V_1^{n-r+1} + \sum_{r=0}^3 C_r^{22} V_2^{n-r+1} - \sum_{r=1}^3 d_r I_2^{n-r+1}. \quad (4)$$

The inductance  $L_1$  should be substituted by  $L_2$  in all the coefficients appearing in Tables I and II of the above paper.

Furthermore, the constants  $q_x$ ,  $\gamma_x$ , and  $\alpha_x(rE_x)$  appearing in (9) of the above paper should be divided by  $d_0$ . Thus, the correct expressions of these equations are, respectively,

$$q_x = 1 + \frac{\Delta t \Delta x C_0^{11}}{2 \varepsilon d_0 \Delta y \Delta z}$$

$$\gamma_x = -\frac{1}{q_x} \frac{\Delta t C_0^{12}}{2 \varepsilon d_0 \Delta y \Delta z}$$

$$\delta_x^n(rE_x) = -\frac{1}{q_x} \left[ \frac{\Delta t}{2 \varepsilon d_0 \Delta y \Delta z} \alpha_x^n(rE_x) - E_x^n(rE_x) - \frac{\Delta t}{\varepsilon} [\nabla \times H]_x^{n+(1/2)}(rE_x) - \frac{\Delta t}{2 \varepsilon} J_{tx}^n(rE_x) \right]. \quad (5)$$

The last error encountered in the above paper concerns two terms of (14), which should be multiplied by  $d_0$ . Thus, this equation should be written as follows:

$$2d_0 f(V_2^{n+1}, V_2^n) + [C_0^{21} \Delta x \gamma_x + C_0^{22}] V_2^{n+1} + [C_0^{21} \Delta x \delta_x^n(rE_x) + \beta_x^n(rE_x) + d_0 I_2^n] = 0. \quad (6)$$